

Down-scaling of resistive switching to nanoscale using porous anodic alumina membranes†

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An advanced approach for resistive switching memory cells based on porous anodic alumina (Al₂O₃) membranes is reported. The effective resistive switching resulting in 6 orders of magnitude difference in resistivity between "on" and "off" states of the cell is achieved by specific electronic and ionic interaction between Ag nanowires filled in the membrane and an ionic conductor (Ag_xAsS₂) deposited on the membrane by thermal evaporation. This easy and robust approach can be exploited for deposition of other ionic conductors for novel types of memories.

Introduction

Current electronic devices (including mobile phones, tablets, and cameras) require increasingly more powerful batteries and memories for their operation. Based on the market demand, there has been for several past years a great attention given to the exploration of new data storage approaches that would enable an increase in the capacity and the speed of memories. Nowadays, conventional data storage technologies, based on charge storage (*e.g.*, DRAM, Flash), are close to their limits.¹ Therefore, there is an urgent need for development of new memory devices based on novel concepts. Several feasible data storage approaches have recently been explored, such as spin-based,² phase change³ and resistive switching⁴ memories. The resistive switching memories, whose differences between the resistivity states enable storage of information, are based on several different mechanisms: thermal, electronic and ionic effects.⁴ In particular, cells based on the ionic effect currently offer intriguing possibilities for more sophisticated data storage applications, such as memristors⁵ or brain like computation.⁶ Moreover, recently published data on these memories seem to be very promising: switching time shorter than 35 ns,⁷ smooth operation for more than 10¹¹ cycles,⁸ and very low energy consumption.⁹

Such performances stem essentially from nanoscale operations. Thus, memories and their parts – cells – have to feature essentially all their dimensions on the nanoscale. Although it is already clear that scale reduction of all memory cell constituents towards the nanoscale is a step forward for memories with

improved performances, the exact switching mechanism in the cell is not yet well understood and has to be further investigated. This is in particular true for memories and their cells based on the ionic switching, where either cations⁸ or anions¹⁰ play a crucial role in the switching process. Some hurdles in understanding the mechanism may arise from the fact that obtaining very small cell parts with a very high precision, closely confined and contamination-free, is not an easy task. As a matter of these facts, thin-film deposition techniques have been employed to prepare these cells. Thermal evaporation,⁷ magnetron sputtering,¹¹ or laser ablation¹² can create homogeneous thin films, whose thickness can be controlled. On the other hand the control of other two dimensions (width and length) remains difficult. As alternatives, common lithographical techniques, such as photolithography,¹³ UV nano-imprint lithography¹⁴ and electron beam lithography,⁷ have been employed. However, these techniques have also their limits in the resolution of the features and may lead to mechanical delamination of cell layers and contamination from the solvents, resins, *etc.*

In this work we demonstrate an advanced approach for the construction of memory cells based on porous anodic alumina membranes (further denoted as AAO) filled with Ag nanowires and covered by a layer of ionic conductor (amorphous chalcogenide glass), where essentially a number of nanowires interfaced with the ionic conductor is responsible for the switching performance providing a sum of individual local contributions of their electrical fields upon switching. In this approach, the AAO membrane with long channels (μm scale) created by small pores (nm scale) acts as a matrix for a large number of electrodeposited Ag nanowires filled homogeneously over the membrane. After the ionic conductor has been deposited on the top of such AAO membrane, the as-filled Ag nanowires become connected with it. Thus, at their interface many small memory electrodes are formed that are actively involved in the switching process and have specific switching characteristics. This very

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effective, yet simple approach is significantly different from memory cells based on stacked multilayers,¹¹ because here instead of forming purely layered electrode systems, we create a huge number of local nanosized electrodes. In general, AAO membranes have offered intriguing possibilities for synthesis of various template-assisted nanowires of metals,^{15,16} polymers,¹⁷ semiconductors¹⁸ and advanced nanostructured systems.^{19,20} In a previous study of one of the coauthors (K. Terabe), AAO membranes have been used for construction of crystalline Ag₂S based resistive memory cells, which were prepared by sulphurization of the Ag layer.²¹ However, in that work the thickness of Ag₂S was approximately 5 μm, which is too thick to explore nanoscale interfaces used in the present work. There have been other studies focusing on the use of single nanowire switching cells²² and exploration of various metal nanowires often with even a narrower diameter than in this work.^{23,24} However, the exploration of thin layers of an amorphous ionic conductor present on an AAO membrane as key components of the resistive memory cells has never been reported. Compared to previous literature,²¹ this approach can be used for a range of ionic conductors. As a case example, we selected amorphous chalcogenide As₂S₂, since it can readily accept Ag ions by photoinduced diffusion and dissolution, and electromigration and thus a variety of compositions can be prepared.²⁵

For clarity of this paper and in particular for the discussion of the results, we want to firstly introduce and briefly discuss the working principle of the memory cell based on the cationic resistive switching, where the metal dissolution in the ionic conductor is the key process in this memory type.^{21,26} As shown in Fig. 1a, the cell consists of distinct layers of diffusive metal (Ag), ion conducting material (Ag-photodoped As₂S₂) and non-diffusive metal as a contact (Al). When a positive bias is applied on the diffusive metal (as depicted in Fig. 1b), the metal undergoes oxidation and its cations are dissolved in the ion conducting material, through which they further migrate towards the cathode (point I). At a certain breakpoint the cations are reduced on the cathode (made of Al) in such a quantity that they form a metallic filament through the ion conducting material (point II). Thus both the electrodes become connected by the metallic filament and the cell resistivity rapidly decreases. On the other hand, applying a negative bias causes dissolution of the conducting filament, while the overall resistivity remains unchanged for some time. Right after

another breakpoint (point III) the overall resistivity of the cell increases again, till its original value (point IV). This can essentially be repeated for many times, until the ionic conductor layer is capable of forming/dissolving the filament accompanied by large resistivity difference. Typically, the cell lifetime is limited by the rupture of the layered cell structure, which may be caused by the heat generation and/or mass flow. Therefore, the potential window has to be carefully selected to maintain flowing current at such a level to maintain long cell lifetime (represented by a large number of cycles without damage). The discussion about the filament built-up will be provided later in the Results and discussion section.

Experimental

Fig. 2 shows in detail two types of memory cells used in the present study. Their electrochemical switching processes follow all particular stages in the switching scheme (Fig. 1), but their interlayer construction is different. The first cell type (Fig. 2a) is based on the AAO membrane filled with Ag nanowires (further denoted as CELL AAO). For these samples, we used commercially available 50 μm-thick Whatman AAO membranes with open pores on both sides (with an average pore diameter of approximately 220 nm). The surface roughness of such membranes is up to 1 μm due to grains and their boundaries. Subsequently, we covered always one side of the membrane by an evaporated Ag layer (300 nm thick) that served as a cathode during Ag electrochemical infilling of AAO. The filling of membranes stems from procedures described in the literature.²⁷ In brief, pulse deposition was carried out in electrolytes consisting of 45 g l⁻¹ AgNO₃, 77 g l⁻¹ CH₃COONH₄, 32 ml l⁻¹ NH₃·H₂O and a small amount of NaOH to adjust the pH to ≈9. The Ag bottom layer of AAO served as the cathode and Pt plate served as the anode. A pulse generator (Agilent 8110A 150 MHz Pulse Generator) was set to repeat 10 ms pulse every 0.5 s. The resulting AAO structures had nearly 100% of the pores filled by Ag. However, their surfaces were covered by Ag overgrown of the pores and some pores were not filled completely (see ESI, Fig. S1†). To remove the overgrown layer and to achieve such depth of AAO membranes, where all pores are filled, the Ag-filled AAO membranes glued to a Cu plate (using Ag-glyce) were ground using SiC emery paper (with grid 2500) down to

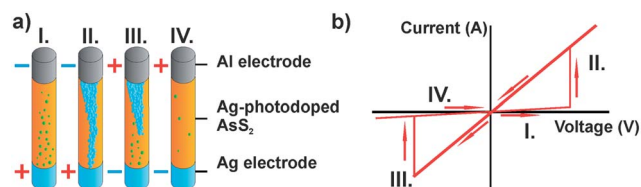


Fig. 1 (a) The scheme of operation of a memory cell based on ionic resistive switching, thicknesses of all layers are on the order of hundreds of nanometers. Symbols + and – represent applied polarities during each particular stage. Green and blue points represent Ag⁺ and Ag⁰ species, respectively. (b) The scheme of 4 stages through which the cell passes within each cycle of operation.

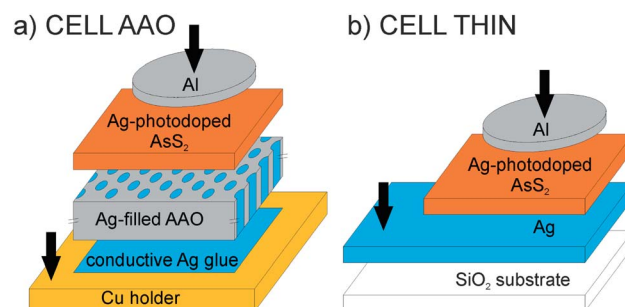


Fig. 2 (a) Scheme of CELL AAO and (b) scheme of CELL THIN. Arrows indicate contacting points.



25 μm and polished in order to achieve a smooth surface (roughness ≈ 10 nm), perpendicular to the direction of Ag nanowires.

The second cell type is based on multilayers of thin films of same components, but without the AAO membrane. For these cells (in further text denoted as CELL THIN), we used silica glass as the starting substrate that was coated by Ag through evaporation resulting in an approx. 300 nm thick layer.

Next steps were identical for both cell types. In brief, both structures were covered by approx. 100 nm thin films of AsS_2 using the thermal evaporation technique. Insulating AsS_2 was modified into ion conducting Ag_xAsS_2 by photo-induced diffusion of Ag. At the end, Al non-diffusive electrodes (approx. 200 nm in thickness, 1 mm^2 surface area) were evaporated on cell surfaces by thermal evaporation of Al through an appropriate mask. Amorphous AsS_2 chalcogenide glass²⁸ was selected as a model ionic conductor from the great family of ionic conductors, including oxides,¹⁰ polymers,²⁹ amorphous³⁰ and crystalline chalcogenides.³¹ Amorphous glassy chalcogenides feature themselves by simple preparation in bulk and also in thin film form,³² and by tuneable properties as ion/electron conductivity ratio,²⁸ content of chalcogen,³³ chemical stability³⁴ and phase separation.³⁵ In contrast to the previous work,²¹ the ionic conductor was purely amorphous here. Since AsS_2 itself is a very good insulator, it was photodoped by Ag (ref. 36) in order to increase its ionic conductivity by light emitted from a mercury lamp for 30 minutes under pure N_2 (purity 4 N) atmosphere. Ag electrodes served as the source of Ag for this photo-diffusion process.

Finally, Al electrodes were deposited by thermal evaporation of Al on the Ag-photodoped AsS_2 . A metallic (Al) mask was used in order to achieve circular electrodes with a diameter of 1 mm^2 . All necessary steps were optimized and the real structure after each step was carefully observed using a scanning electron microscope (Mira3 XM FEG SEM, Tescan). All in all, these procedures resulted in two types of principally very similar cell types, but having strongly different design and dimensions of the Ag diffusive electrode.

Results and discussion

Fig. 3 shows SEM images of an empty AAO membrane and an Ag-filled AAO membrane as an example of CELL AAO. For comparison, we provide top-views of the empty membrane (Fig. 3a) and the polished Ag-filled AAO membrane (Fig. 3b). From these images one can see that after polishing the surface becomes very uniform and smooth. Moreover, from Fig. 3c and d that display the whole cell in the cross-sectional view, it becomes clear that the original membrane thickness of 50 μm was ground down to approximately 25 μm and polished, resulting in a smooth membrane having continuous embedded nanowires in the pores. The membrane thinning to approximately half of its original thickness is responsible for the fact that the original pore diameter at the membrane top decreased from approximately 220 nm to approximately 190 nm. Moreover, Fig. 3e and f show the same layers in cross-sectional views under higher magnifications, where contrast difference

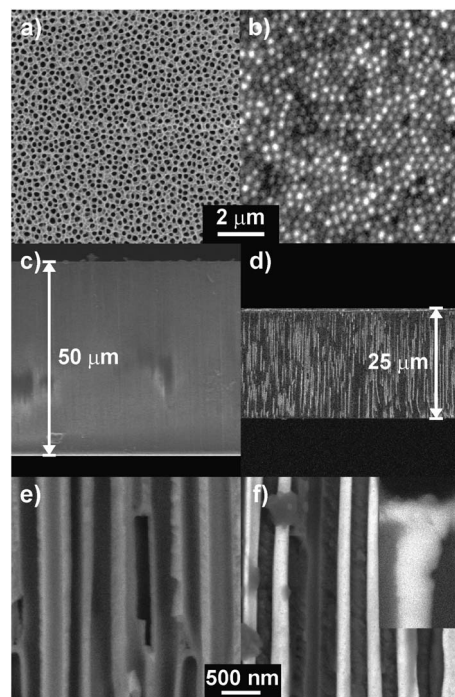


Fig. 3 SEM images of CELL AAO. Left column displays images of an empty AAO membrane, right column displays images of a polished Ag-filled AAO membrane. We show top-views (a and b), cross-sectional views of the membrane (c and d) and high magnifications of the same (e and f). The inset of (f) shows the interface between Ag-filled AAO and Ag_xAsS_2 after switching experiments have been performed (10 cycles).

between AAO and Ag (owing to different quantities of back-scattered electrons) reveals homogeneous filling of AAO pores by Ag nanowires. Finally, the inset of Fig. 3f shows a detailed cross-sectional view of the interface between Ag-filled AAO and Ag_xAsS_2 after switching experiments have been performed (10 cycles).

As one can see from the white appearance (employing signal from the back-scattered electron detector), resulting Ag_xAsS_2 contains quite a significant amount of Ag. However, its exact amount cannot be measured accurately by surface analytical techniques (such as EDX or XPS), because these techniques enhance Ag migration.

Cell designs used in the present study allow us to investigate the effect of the size and shape of the Ag electrode on the resistive switching behavior. Therefore, the as-prepared samples were subjected to cyclic voltammetry. Fig. 4 shows comparison of voltammograms recorded for samples of CELL AAO (Fig. 4a) and CELL THIN (Fig. 4b). A detailed description of these measurements is given in the ESI.† For the purpose of clarity, we provide response in terms of current (left axes) as well as current densities (right axes). Both plots show data recorded within four measurement cycles: 1st, 5th, 10th and 20th, all of them reflect typical switching behavior as described in Fig. 1. For both cell types, the switch “on” and switch “off” occur abruptly and the threshold voltages are below 0.3 V and −0.3 V, respectively. The obvious differences in current levels and shifts



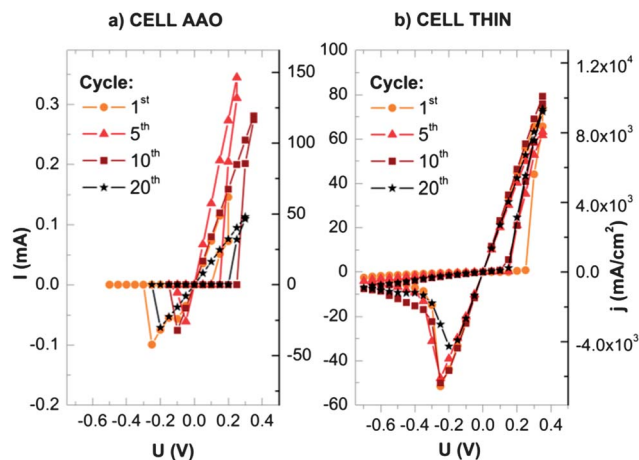


Fig. 4 Cyclic voltammograms of (a) CELL AAO and (b) CELL THIN for several particular cycles. Left Y-axes display currents and right Y-axes display current densities.

of curves upon cycling stem partially from the difference in contact areas on samples during switching, but mainly from the changing stoichiometry. The contact area for CELL AAO is approximately 30% of the contact area of the CELL THIN (as it was calculated from the porosity of AAO assuming that all its pores are filled with Ag, which is almost the case after polishing as demonstrated in Fig. 3). The differences in performance shown in Fig. 4 originating from changing stoichiometry will be discussed later. The data presented in Fig. 4 are sufficient for the demonstration of the switching of the cells. However, they do not provide straightforward information about cyclability and differences in resistivities upon extended cycling.

Therefore, recorded voltammetric data were further analyzed (details are given in the ESI†) in order to point out the resistivity of the “on” and “off” states in each cycle. Results are plotted in Fig. 5 that shows the “on” and “off” resistivity for both cell types and their evolution upon cycling. In addition resistivities of studied cells before the first switch “on” are also presented, so the overall history of the sample is recorded.

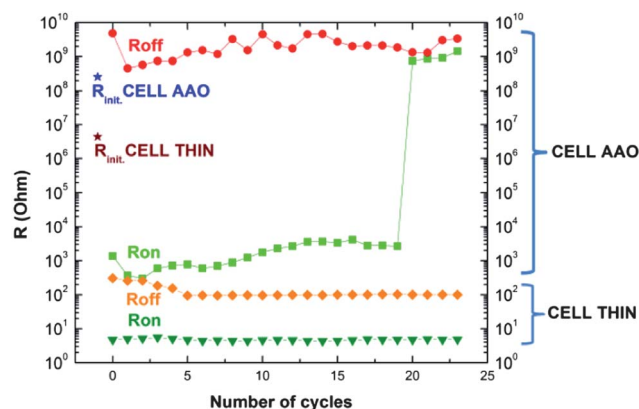


Fig. 5 Comparison of resistance data between the CELL AAO and the CELL THIN.

From Fig. 5 a huge change in “on” and “off” state resistivities for CELL AAO in comparison with CELL THIN is clearly visible. In the case of CELL THIN resistivities were much lower (the “on” state resistivity was below $10^1 \Omega$ and the “off” state resistivity was above $10^2 \Omega$) for all measured cycles. It is also noteworthy that its initial resistivity R_{int} was comparably high ($4.4 \times 10^6 \Omega$).

For CELL AAO resistivities were much higher, the “on” state resistivity was below $10^4 \Omega$ and the “off” state resistivity was above $10^8 \Omega$ for all measured cycles. The initial resistivity of the “on” state of the CELL AAO was $2.6 \times 10^8 \Omega$, which is quite close to typical “off” state resistivity. However, in the 21st cycle the resistivity of the “on” state of CELL AAO became quite high and the switching behavior was lost. This behavior was observed throughout all samples, usually after 20 and more cycles.

Taking into account our results and observed findings we can discuss now about the switching mechanism. In CELL THIN the initial resistivity (R_{int}) is much higher than resistivities at the “off” state (R_{off}), *i.e.* immediately within the first switching cycle the Ag stoichiometry in the ionic conductor significantly changes. The lower resistivity indicates that there is a significant uptake of Ag by electromigration into the ionic conductor that leads to a significant decrease of the overall resistivity. The reasons for such an uptake are essentially twofold.

At first, we can assume that the Ag rich area contains conductive Ag⁰ clusters or filaments, as observed experimentally by Guo *et al.*³⁷ and schematically depicted in Fig. 1a. The filament has been shown in the literature to have most likely the dendrite shape.³⁷ The shape and the number of dendrites, however, stem from the amount of dissolved metal in the ionic conductor. The dissolved amount is influenced by the initial metal amount and by the electric field set at a certain potential difference. The filament formation is naturally very desirable and it is the most important phenomenon for the whole cell.⁴ However, at the same time, it is also the most critical factor. If the created filament is too big, it cannot be completely dissolved and the resulting material will not reach its original resistivity.²¹

Secondly, a stoichiometric enrichment of the Ag content in Ag_xAsS₂ may occur that leads to a significant change of the intrinsic resistivity of the Ag_xAsS₂.^{28,38} Naturally, both effects have an impact on decreased resistivity of the cell, no matter whether we speak about the “on” or the “off” state, as their difference is rather low in the case of CELL THIN. It has to be noted that for this cell type, the Ag comes from an almost infinite source of Ag layer and the presence of such rich source must influence the resulting distribution significantly. Without specifying the dominant effect, it is reasonable to assign better reproducibility among switching cycles observed during cycling voltammetry to the fact that the “on” and “off” states are accompanied by relatively minor changes in the layers of CELL THIN in terms of their stoichiometry, not strong enough to alter the electrochemical characteristics during cycling (Fig. 4b).

In contrast, the situation for CELL AAO is different in two points. Firstly, the electric field flowing through the cell is more focused (*i.e.* less randomly oriented), due to the regular alignment of the nanowire Ag-electrode embedded in the AAO matrix. This must have an impact on the creation of conducting pathways through the ionic conductor. Similar effects on



electric field caused by nanosized structures have been described in the literature.^{39,40} The other main point is the limited source of Ag for photoinduced diffusion and filament creation compared to CELL THIN. In other words, there is not as much Ag available in the nanowires as in the bulk Ag layer due to the AAO porosity of 30%. For instance, if we would compare the amount of Ag available for the switching between 300 nm thick Ag layer and 300 nm long Ag nanowires embedded in 300 nm thick AAO membrane, Ag nanowires would account only for 30% of the Ag content compared to the bulk Ag layer. However, the resistance of the “on” and “off” states in the CELL AAO is huge – almost 6 orders of magnitude.

In order to provide a more insight into the CELL AAO interface between the ionic conductor and Ag nanowires embedded in the AAO matrix, we show in Fig. 6 a comparison of images acquired by SEM before the entire switching experiment and after the entire switching experiment with the cell damaged. During the SEM screening, we intentionally used both types of detectors separately (secondary-electron detector, SE, and back-scattered electron, BSE, respectively) and simultaneously to evaluate the connection of Ag nanowires to the layer of ionic conductor in terms of morphology and composition.

As it can be clearly seen from the comparison, the Ag nanowires embedded in the AAO membrane are in solid contact with the ionic conductor before switching (*i.e.* after Ag photodiffusion into the ionic conductor) and naturally during the

switching (as shown in Fig. 3f), while after a critical number of switching cycles (under given conditions used in this work typically between 20 and 25 cycles) they become disconnected as a result of Ag depletion. This is in particular apparent from the BSE image of the switched cell, where a clear gap between the Ag nanowire and the ionic conductor is visible. We have observed these features in all samples involved in our switching experiments. As a matter of fact, the resistive switching is accompanied by relatively strong Ag mass flow (leading to its depletion) on the nanoscale interface between Ag-filled AAO, ionic conductor and Al electrode.

Moreover, this explains the shift between cycles recorded for CELL AAO (Fig. 4a), in particular for the CELL AAO resistance in the “on” state (see R_{on} for CELL AAO in Fig. 5) that shows an overall increase during the switching. This increase essentially provides additional confirmation that upon cycling the resistance of this interface increases until the interface becomes disconnected during one particular switching cycle (as apparent from the R_{on} jump in Fig. 5 for CELL AAO and as corroborated by Fig. 6).

In our view, to describe the features that lead to huge resistivity changes, we have to consider also the concept of conductive filaments as in the case of CELL THIN. However, in this case, a comparably thinner and shorter conductive filament over the ionic conductor is built in the period I shown in Fig. 1a. Nevertheless, it is conductive enough to dramatically decrease resistivity when built (R_{on} level), however easily dissolved in the reverse process with a strong impact on the resistivity (R_{off} level). This concept corresponds to the work of Wang *et al.*²⁶ that assumes that the filament is connected with Ag⁺ rich areas distributed in the ion conductive matrix and thus the volume of the ion conductor between electrodes directly influences the switching behavior. As apparent from Fig. 5, the R_{int} value for CELL AAO is lower than any other value R_{off} recorded afterwards. This can be explained by combination of dissolution of the Ag filament, or lowering the Ag content in non-stoichiometric Ag_xAsS_2 at the Ag– Ag_xAsS_2 interface upon cycling. The continuous migration of Ag from nanowires towards the Al cathode through the ionic electrode results in slow depletion of Ag from Ag nanowires. Ag-depletion of the Ag nanowire continuously increases with each new cycle, till contact is lost (as shown in Fig. 6) and the cell loses its function. Although the Ag-depletion is currently limiting the number of switching cycles that this type of cell is able to achieve, the magnitude of the resistance difference during the switching opens new pathways for an efficient data storage, or large data quantity.

We believe that by careful optimization, the number of achievable cycles could be strongly increased. We foresee several possibilities for a considerable increase of cycle numbers. The easiest option appears to increase the Ag content by deposition of a thin Ag film onto Ag-filled AAO, before depositing the ionic conductor. This Ag layer should have such thickness to become entirely dissolved in the Ag_xAsS_2 during the photodiffusion process. As a result, Ag_xAsS_2 would become Ag enriched (compared to situation without an additional Ag layer) and less Ag from nanowires would be consumed for the switching. Thus the number of achievable switching cycles

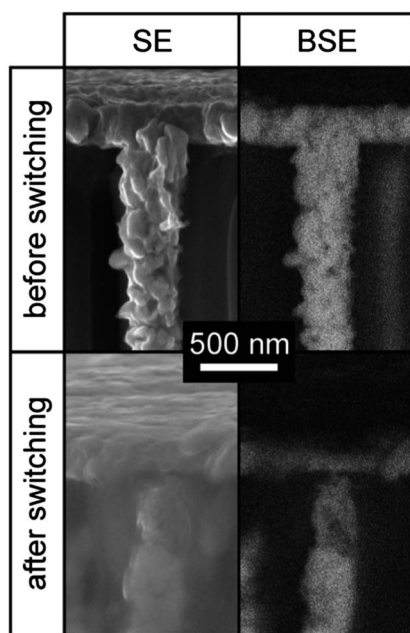


Fig. 6 SEM images of the CELL AAO sample showing the interface of the ionic conductor and Ag nanowires embedded in the AAO matrix. SE stands for imaging using a secondary-electron detector addressing the morphology of the observed samples, BSE stands for a back-scattered electron detector addressing the compositional differences on the interface. Observations were done under a high magnification to see properly details of the interface (magnification 100 000 \times , working distance 5 mm). The sample “after switching” was observed after recording 21 cycles and stopping the experiment afterwards.



would increase, potentially without any loss in the field focusing that helps us to maintain such a distinct resistance change between switching. So we could obtain an ideal compromise between the lifetime and the resistivity difference. It is noteworthy that adding an additional Pt layer in the NiO-based resistive switching cell with a positive effect on the switching has been reported recently.⁴¹ Another possibility could be to improve the electrical contact between Ag nanowires and the chalcogenide ionic conductor by using an even smoother and cleaner Ag-filled AAO surface, prepared by optimized polishing and/or ion milling. It is possible that not all nanowires are electrically connected to the ionic conductor. Eventually, optimized AAO membranes for Ag-filling could be used. This will be a subject of further work.

Conclusions

We demonstrated comparably advanced design of resistive switching memory cells based on the Ag-filled porous anodic alumina membrane. The cells were prepared by low-cost combination of electrodeposition (filling of Ag into the membrane) and the thermal evaporation technique (to deposit layers of an ionic conductor Ag_xAsS_2 and Al contacts). The vital effect of a cell consisting of Ag-filled AAO and an ionic conductor Ag_xAsS_2 on the switching behavior was confirmed by direct comparison of samples with and without an Ag-filled AAO electrode. We achieved differences of 6 orders of magnitude in the resistivities between “on” and “off” states. Considering the great tuneability of dimensions of the AAO membrane and the possibility to use a whole range of ionic conductors, we believe that the presented approach could be explored and optimized for many other novel types of resistive switching memories based on the ionic effect.

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Notes and references

- International Technology Roadmap for Semiconductors: Overall Roadmap Technology Characteristics, Semiconductor Industry Association, Incheon, Korea, 2011, available from <http://public.itrs.net/>.
- C. Chappert, A. Fert and F. N. Van Dau, *Nat. Mater.*, 2007, **6**, 813–823.
- M. D. Ventra and Y. V. Pershin, *Mater. Today*, 2011, **14**, 584–591.
- R. Waser and M. Aono, *Nat. Mater.*, 2007, **6**, 833–840.
- J. J. Yang, D. B. Strukov and D. R. Stewart, *Nat. Nanotechnol.*, 2013, **8**, 13–24.
- T. Hasegawa, T. Ohno, K. Terabe, T. Tsuruoka, T. Nakazama, J. K. Gimzewski and M. Aono, *Adv. Mater.*, 2010, **22**, 1831–1834.
- M. N. Kozicki, M. Mitkova, M. Park, M. Balakrishnan and C. Gopalan, *Superlattices Microstruct.*, 2003, **34**, 459–465.
- M. N. Kozicki, M. Park and M. Mitkova, *IEEE Trans. Nanotechnol.*, 2005, **4**, 331–338.
- C. Schindler, M. Meier, R. Waser and M. N. Kozicki, *Non-volatile Memory Technology Symposium*, 2007, vol. 82–85, pp. 10–13.
- R. Yang, K. Terabe, T. Tsuruoka, T. Hasegawa and M. Aono, *Appl. Phys. Lett.*, 2012, **100**, 231603.
- C. Schindler, I. Valov and R. Waser, *Phys. Chem. Chem. Phys.*, 2009, **11**, 5974–5979.
- X. B. Yan, J. Yin, H. X. Guo, Y. Su, B. Xu, H. T. Li, D. W. Yan, Y. D. Xia and Z. G. Liu, *J. Appl. Phys.*, 2009, **106**, 054501.
- J. van den Hurk, I. Valov and R. Waser, *Thin Solid Films*, 2013, **527**, 299–302.
- M. Meier, S. Gilles, R. Rosezin, C. Schindler, S. Trelenkamp, A. Rüdiger, D. Mayer, C. Kügeler and R. Waser, *Microelectron. Eng.*, 2009, **86**, 1060–1062.
- K. Nielsch, F. Müller, A.-P. Li and U. Gösele, *Adv. Mater.*, 2000, **12**, 582–586.
- J. M. Baik, M. Schierhorn and M. Moskovits, *J. Phys. Chem. C.*, 2008, **112**, 2252–2255.
- T.-H. Fang, T. H. Wang and S.-H. Kang, *J. Mater. Sci.*, 2009, **44**, 1588–1593.
- W.-B. Zhao, J.-J. Zhu and H.-Y. Chen, *Scr. Mater.*, 2004, **50**, 1169–1173.
- W.-C. Yoo and J.-K. Lee, *Adv. Mater.*, 2004, **16**, 1097–1101.
- S.-H. Lyu and J.-S. Lee, *J. Mater. Chem.*, 2012, **22**, 1852–1861.
- C. Liang, K. Terabe, T. Hasegawa, R. Negishi, T. Tamura and M. Aono, *Small*, 2005, **1**, 971–975.
- C. Liang, K. Terabe, T. Hasegawa and M. Aono, *Nanotechnology*, 2007, **18**, 485202.
- D. Perego, S. Franz, M. Bestetti, L. Cattaneo, S. Brivio, G. Tallarida and S. Spiga, *Nanotechnology*, 2013, **24**, 045302.
- D. Ielmini, C. Cagli, F. Nardi and Y. Zhang, *J. Phys. D: Appl. Phys.*, 2013, **46**, 074006.
- T. Wagner, S. O. Kasap, Mir. Vlcek, M. Frumar, P. Nesladek and Mil. Vlcek, *Appl. Surf. Sci.*, 2001, **175–176**, 117–122.
- F. Wang, W. P. Dunn, M. Jain, C. De Leo and N. Vickers, *Solid-State Electron.*, 2011, **61**, 33–37.
- G. Sauer, G. Brehm, S. Schneider, K. Nielsch, R. B. Wehrspohn, J. Choi, H. Hofmeister and U. Gösele, *J. Appl. Phys.*, 2002, **91**, 3243–3247.
- S. Stehlik, J. Kolar, M. Bartos, M. Vlcek, M. Frumar, V. Zima and T. Wagner, *Solid State Ionics*, 2010, **181**, 1625–1630.
- S. Wu, T. Tsuruoka, K. Terabe, T. Hasegawa, J. P. Hill, K. Ariga and M. Aono, *Adv. Funct. Mater.*, 2011, **21**, 93–99.
- I. Valov, R. Waser, J. R. Jameson and M. N. Kozicki, *Nanotechnology*, 2011, **22**, 254003.
- K. Terabe, T. Hasegawa, T. Nakayama and M. Aono, *Nature*, 2005, **433**, 47–50.
- A. B. Seddon, *J. Non-Cryst. Solids*, 1995, **184**, 44–50.
- S. N. Yannopoulos, F. Kyriazis and I. P. Chochliouros, *Opt. Lett.*, 2011, **36**, 534–536.



- 34 J. Orava, T. Wagner, M. Krbal, T. Kohoutek, Mil. Vlcek and M. Frumar, *J. Non-Cryst. Solids*, 2006, **352**, 1637–1640.
- 35 S. Stehlik, J. Kolar, H. Haneda, I. Sakaguchi, M. Frumar and T. Wagner, *Int. J. Appl. Glass Sci.*, 2011, **2**, 301–307.
- 36 M. Frumar and T. Wagner, *Curr. Opin. Solid State Mater. Sci.*, 2003, **7**, 117–126.
- 37 X. Guo, C. Scindler, S. Menzel and R. Waser, *Appl. Phys. Lett.*, 2007, **91**, 133513.
- 38 M. Ohto, M. Itoh and K. Tanaka, *J. Appl. Phys.*, 1995, **77**, 1034–1039.
- 39 Q. Liu, S. Long, W. Wang, S. Tanachutiwat, Y. Li, Q. Wang, M. Zhang, Z. Huo, J. Chen and M. Liu, *IEEE Electron Device Lett.*, 2010, **31**, 1299–1301.
- 40 Z. Q. Wang, H. Y. Xu, L. Zhang, X. H. Li, J. G. Ma, X. T. Zhang and Y. C. Liu, *Nanoscale*, 2013, **5**, 4490–4494.
- 41 Y.-C. Huang, P.-Y. Chen, T.-S. Chin, R.-S. Liu, C.-Y. Huang and C.-H. Lai, *Appl. Phys. Lett.*, 2012, **101**, 153106.

